

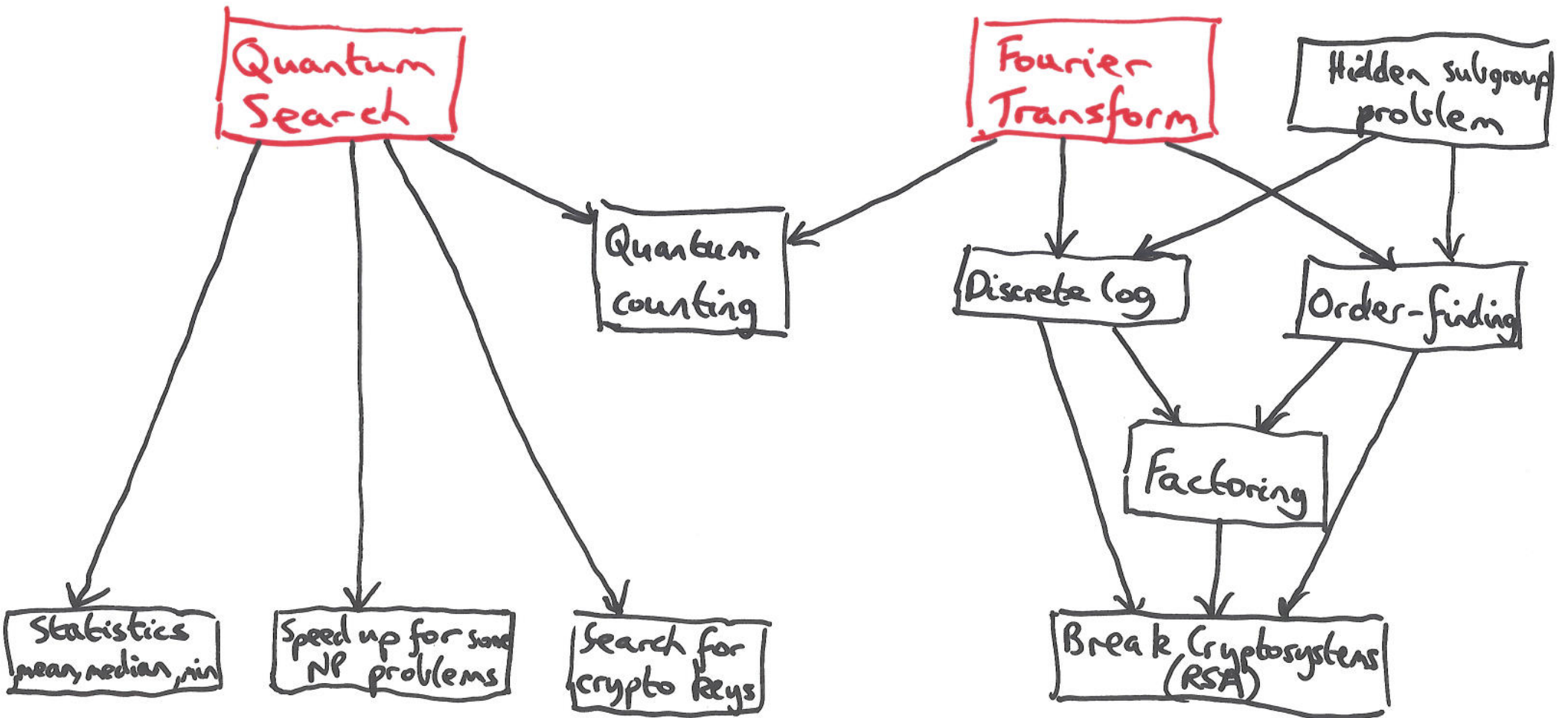
# Quantum Circuits

## Quantum Algorithms

Two classes:

1. Based on **Shor's quantum Fourier transform**
  - exponential speed-up.
2. Based on **Grover's algorithm for quantum searching**
  - quadratic speed-up.

# Main quantum algorithms



Why so few quantum algorithms?

... a difficult problem.

Why?

1. Algorithm design is difficult  
... must be better than best classical.
2. Our intuitions adapted to classical world.

# Single Qubit Operations

e.g.

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}; \quad Y = i \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}; \quad Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}; \quad S = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}; \quad T = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$$

← Sometimes called  $\frac{\pi}{8}$  gate for silly reasons

Observe,

$$H = (X+Z)/\sqrt{2}, \quad S = T^2$$

A single qubit,  $a|0\rangle + b|1\rangle$  can be visualized as  $(\theta, \varphi)$  on the unit sphere, where  $a = \cos(\theta/2)$ ,  $b = e^{i\varphi} \sin(\theta/2)$ .  
a can be considered real up to overall phase.

Bloch vector is  $(\cos \varphi \sin \theta, \sin \varphi \sin \theta, \cos \theta)$

## Exercise

Find the points on the Bloch sphere which correspond to the normalized eigenvectors of the different Pauli matrices.

e.g. eigenvectors of  $X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$  are  $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$  and  $\begin{pmatrix} 1 \\ -1 \end{pmatrix}$

$\Rightarrow$  normalised eigenvectors are  $(|0\rangle + |1\rangle)/\sqrt{2}$  and  $(|0\rangle - |1\rangle)/\sqrt{2}$

For  $(|0\rangle + |1\rangle)/\sqrt{2}$ :

$$a|0\rangle + b|1\rangle \Rightarrow a = b = \frac{1}{\sqrt{2}}$$

$$\text{But } a = \cos(\theta/2) \Rightarrow \theta = \frac{\pi}{2}, \quad b = e^{i\varphi} \sin(\theta/2) \Rightarrow \varphi = 0$$

$$\Rightarrow \text{Bloch vector} = (\cos\varphi \sin\theta, \sin\varphi \sin\theta, \cos\theta) = (1, 0, 0).$$



# Exponents of Pauli matrices

Rotation operators about the  $\hat{x}$ ,  $\hat{y}$  and  $\hat{z}$  axes,

$$R_x(\theta) = e^{-i\theta X/2} = \cos\frac{\theta}{2} I - i\sin\frac{\theta}{2} X = \begin{bmatrix} \cos\frac{\theta}{2} & -i\sin\frac{\theta}{2} \\ -i\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$

$$R_y(\theta) = e^{-i\theta Y/2} = \cos\frac{\theta}{2} I - i\sin\frac{\theta}{2} Y = \begin{bmatrix} \cos\frac{\theta}{2} & -\sin\frac{\theta}{2} \\ \sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$

$$R_z(\theta) = e^{-i\theta Z/2} = \cos\frac{\theta}{2} I - i\sin\frac{\theta}{2} Z = \begin{bmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{bmatrix}.$$

Let  $x$  be a real number and  $A$  a matrix such that  $A^2 = I$ . Then,

$$e^{iAx} = \cos(x)I + i\sin(x)A$$

Proof sketch: eigenvalues of  $A$  are  $\pm 1$ . e.g. if both  $\lambda_0 = \lambda_1 = \pm 1$ , then

$$\begin{aligned} e^{iAx} &= e^{i\lambda_0 x} |a_0\rangle\langle a_0| + e^{i\lambda_1 x} |a_1\rangle\langle a_1| \\ &= \cos(x) [ |a_0\rangle\langle a_0| + |a_1\rangle\langle a_1| ] + i\sin(x) [ \lambda_0 |a_0\rangle\langle a_0| + \lambda_1 |a_1\rangle\langle a_1| ] \\ &= \cos(x)I + i\sin(x)A. \end{aligned}$$

$$\Rightarrow \text{e.g. } R_x(\theta) = e^{-i\theta X/2} = \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}X.$$

Similarly,

$$T = R_3\left(\frac{\pi}{4}\right) = \begin{bmatrix} e^{-i\frac{\pi}{2}} & 0 \\ 0 & e^{i\frac{\pi}{2}} \end{bmatrix},$$

to within  
global phase.

Also,

$$H = \begin{bmatrix} \cos \frac{\pi}{4} & -\sin \frac{\pi}{4} \\ \sin \frac{\pi}{4} & \cos \frac{\pi}{4} \end{bmatrix} \begin{bmatrix} e^{-i\pi} & 0 \\ 0 & e^{i\pi} \end{bmatrix} \times^{-1} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$= R_x\left(\frac{\pi}{2}\right) \times R_3(2\pi) \times e^{i\pi}.$$



More generally,

$$R_{\hat{n}}(\theta) = e^{-i\theta \hat{n} \cdot \vec{\sigma} / 2} = \cos\left(\frac{\theta}{2}\right) \mathbb{I} - i \sin\left(\frac{\theta}{2}\right) (n_x X + n_y Y + n_z Z),$$

where  $\hat{n} = (n_x, n_y, n_z)$ , and  $\vec{\sigma} = (X, Y, Z)$ .

This can be proved by using,

$$(\hat{n} \cdot \vec{\sigma})^2 = \mathbb{I}.$$

# Bloch Sphere interpretation

Let a qubit be represented by  
Bloch vector,  $\vec{\lambda}$ .

Then  $R_{\hat{n}}(\theta)$  rotates the state by  
angle  $\theta$  about the  $\hat{n}$  axis.

Observe:

$$X R_y(\theta) X = R_y(-\theta).$$

Proof:

$$\text{Use } XYX = -Y.$$

# Single Qubit Unitary

$$U = e^{i\alpha} R_{\hat{n}}(\theta),$$

Prove?

$$H = X + Z$$

$$\Rightarrow H = e^{i\frac{\pi}{2}} R_{101}(\pi).$$

$$S = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} = e^{i\frac{\pi}{4}} \begin{bmatrix} e^{-i\frac{\pi}{4}} & 0 \\ 0 & e^{i\frac{\pi}{4}} \end{bmatrix} = e^{i\frac{\pi}{4}} R_{001}\left(\frac{\pi}{2}\right).$$

## Z-Y Decomposition for a Single Qubit

$$U = e^{i\alpha} R_z(\beta) R_y(\gamma) R_z(\delta).$$

Proof:

$$U = \begin{bmatrix} e^{i(\alpha - \beta/2 - \delta/2)} \cos \frac{\gamma}{2} & -e^{i(\alpha - \beta/2 + \delta/2)} \sin \frac{\gamma}{2} \\ e^{i(\alpha + \beta/2 - \delta/2)} \sin \frac{\gamma}{2} & e^{i(\alpha + \beta/2 + \delta/2)} \cos \frac{\gamma}{2} \end{bmatrix}.$$



Remember:

$$U = e^{i\alpha} R_3(\beta) R_y(\gamma) R_3(\delta).$$

Give a decomposition using  $R_x$  instead of  $R_3$

.....

## Another Decomposition

$$U = e^{i\alpha} R_{\hat{n}}(\beta) R_{\hat{m}}(\gamma) R_{\hat{n}}(\delta),$$

where  $\hat{n}$  and  $\hat{m}$  are not parallel.

## Another Decomposition

$\exists$  unitaries  $A, B, C$ , such that,

$$ABC = I, \text{ and}$$

$$U = e^{i\alpha} A X B X C.$$

Proof:

$$\text{Set } A = R_3(\beta) R_y(\gamma/2), \quad B = R_y(-\gamma/2) R_3(-(\delta+\beta)/2), \\ C = R_3((\delta-\beta)/2).$$

$$\text{Then, } ABC = I, \quad X B X = R_y(\gamma/2) R_3(\frac{\delta+\beta}{2}), \quad A X B X C = R_3(\beta) R_y(\gamma) R_3(\delta)$$

$$U = e^{i\alpha} AXC$$

What are A, B, and C, and  $\alpha$  for the Hadamard gate?

Remember:

$$U = e^{i\alpha} R_z(\beta) R_y(\gamma) R_z(\delta)$$

$$= e^{i\alpha} AXC$$

$$= \begin{bmatrix} e^{i(\alpha - \beta/2 - \delta/2)} \cos \gamma/2 & -e^{i(\alpha - \beta/2 + \delta/2)} \sin \gamma/2 \\ e^{i(\alpha + \beta/2 - \delta/2)} \sin \gamma/2 & e^{i(\alpha + \beta/2 + \delta/2)} \cos \gamma/2 \end{bmatrix}$$

...

$$\dots = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \dots ??$$

Observe:

$$HXH = Z, \quad HYH = -Y, \quad HZH = X,$$

$$HTH = R_x(\pi/4), \quad \dots \text{to within global phase.}$$

... therefore  $X$  and  $Z$  are related by conjugacy with respect to  $H$ ,  
and  $Y$  is self-conjugate with respect to  $H$ .



# Composition of Single Qubit Operations

Rotate by  $\beta_1$  about  $\hat{n}_1$ .

then Rotate by  $\beta_2$  about  $\hat{n}_2$ .

Result is rotation by  $\beta_{12}$  about  $\hat{n}_{12}$ ,  
where,

$$c_{12} = c_1 c_2 - s_1 s_2 \hat{n}_1 \cdot \hat{n}_2$$

$$s_{12} \hat{n}_{12} = s_1 c_2 \hat{n}_1 + c_1 s_2 \hat{n}_2 - s_1 s_2 \hat{n}_2 \times \hat{n}_1,$$

where  $c_i = \cos(\beta_i/2)$ ,  $s_i = \sin(\beta_i/2)$ ,

$$c_{12} = \cos(\beta_{12}/2), \quad s_{12} = \sin(\beta_{12}/2).$$

If  $\beta_1 = \beta_2$  and  $\hat{n}_1 = \hat{z}$ , then,

$$c_{12} = c^2 - s^2 \hat{z} \cdot \hat{n}_2, \quad s_{12} \hat{n}_{12} = s c (\hat{z} + \hat{n}_2) - s^2 \hat{n}_2 \times \hat{z},$$

where  $c = c_1$ , and  $s = s_1$ .

# Common Circuit Symbols

$$\boxed{H} \quad \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$\boxed{X} \quad \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\boxed{Y} \quad \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

$$\boxed{Z} \quad \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

$$\boxed{S} \quad \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

$$\boxed{T} \quad \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$$

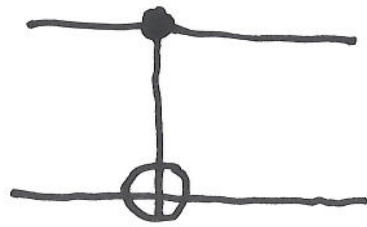
# Controlled Operations

" If A is true,  
then do B".

# Controlled-NOT (CNOT)

$$|c\rangle|t\rangle \rightarrow |c\rangle|t \oplus c\rangle$$

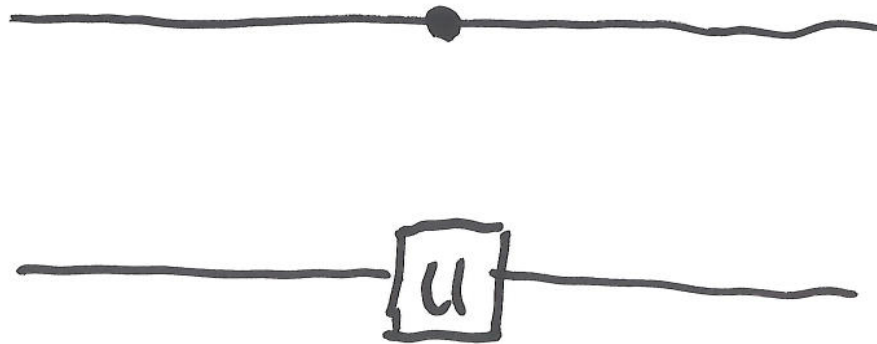
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$



More generally . . . .

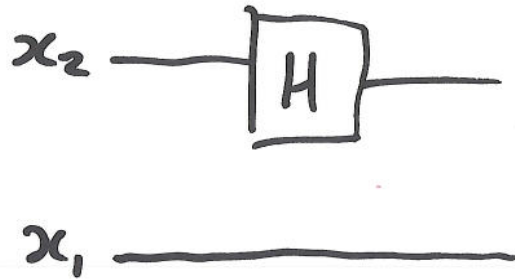
Controlled-U

$$|c\rangle|t\rangle \rightarrow |c\rangle U^c|t\rangle$$



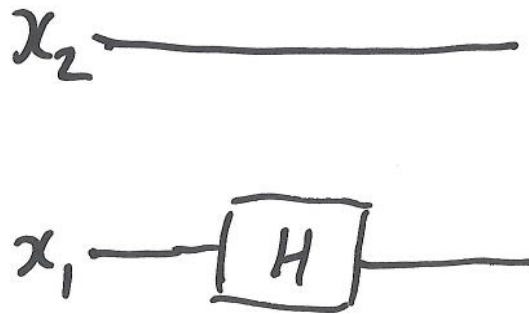


What is the  $4 \times 4$  unitary matrix for



in the computational basis?

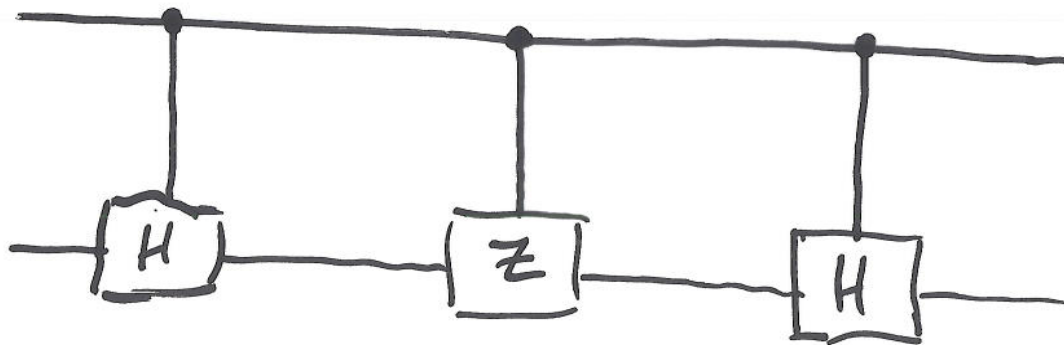
and for



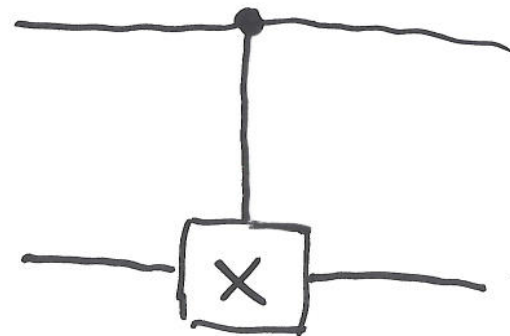
?

# Controlled-Z Gate

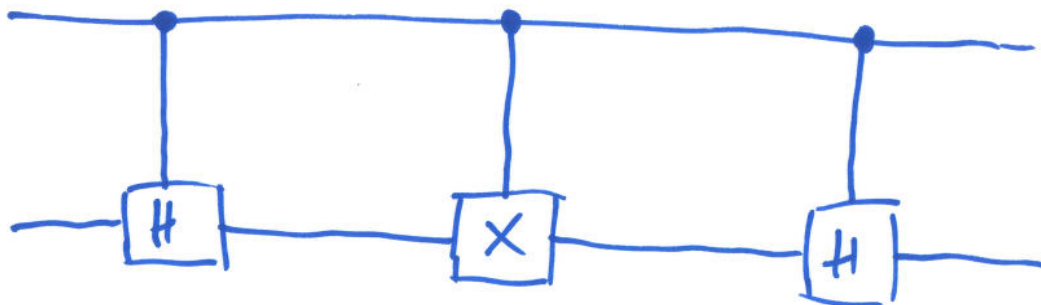
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ 0 & 0 & \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}
 \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}
 \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ 0 & 0 & \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}
 =
 \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$



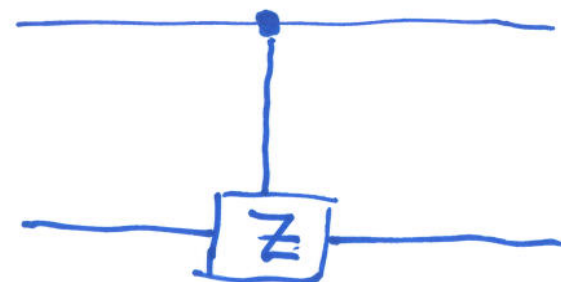
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Similarly

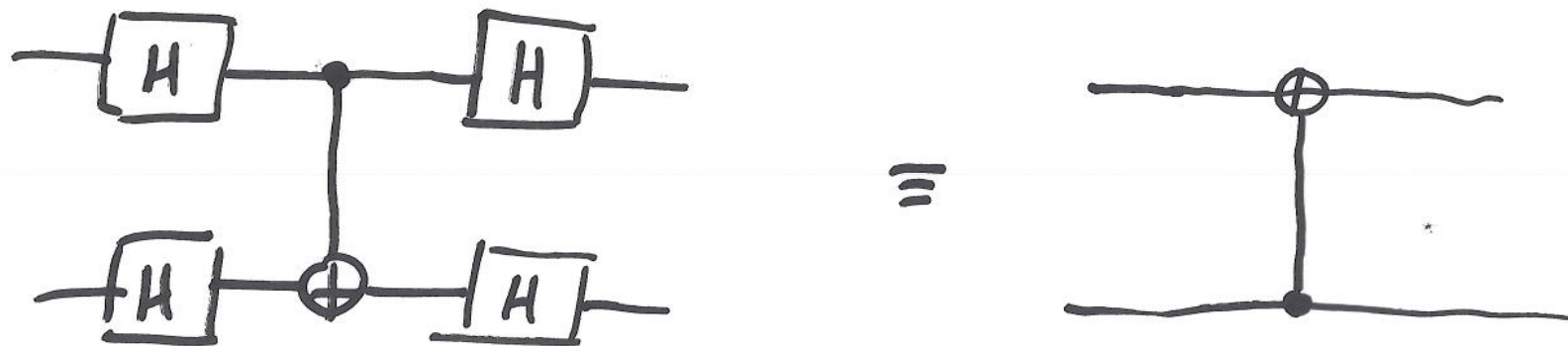


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# Basis Transformations

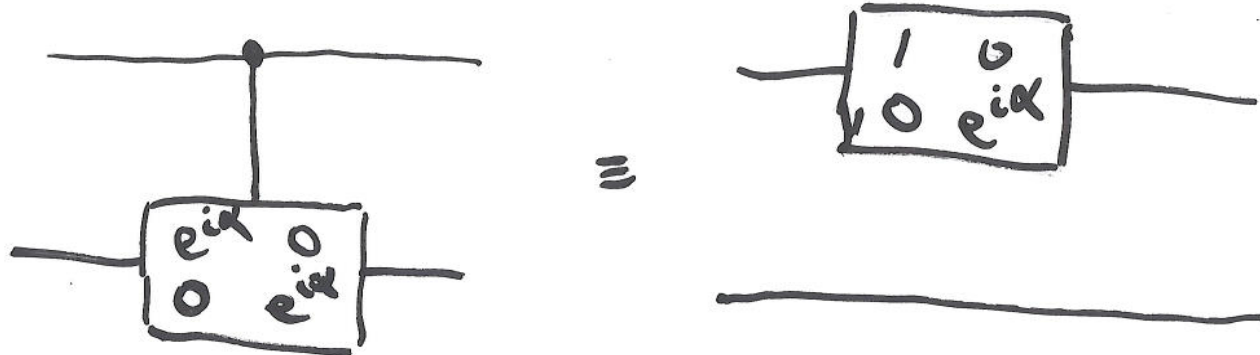
# Conjugacy Relationships



Show this.

Implement Controlled-U using single qubit operations and controlled-NOT

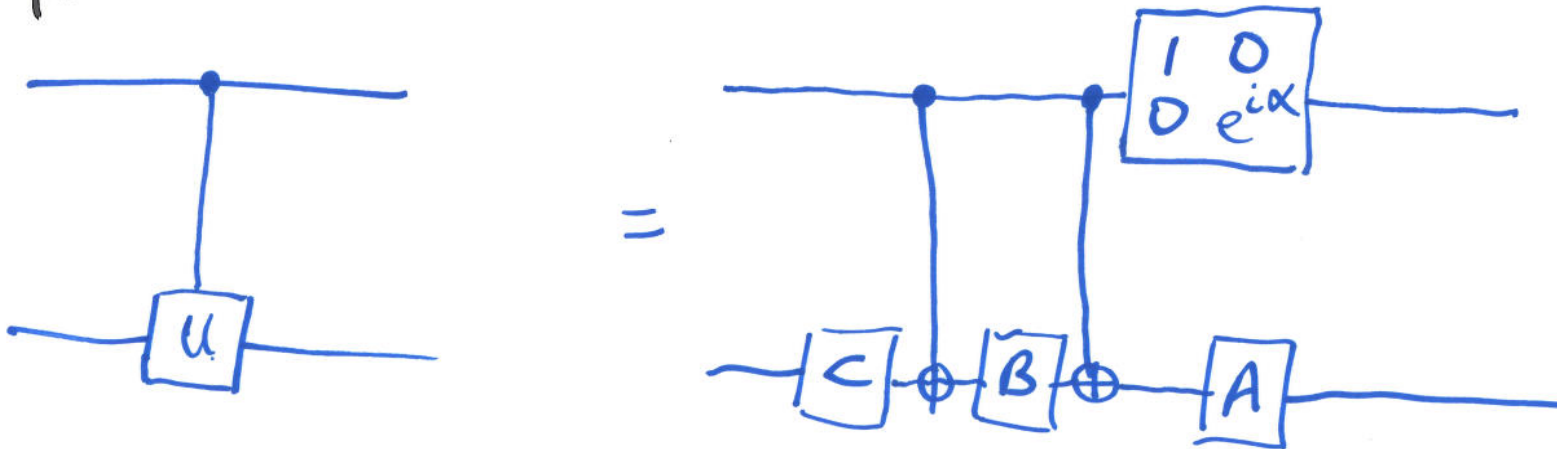
observe:



also,

$$U = e^{i\alpha} AXC$$

therefore

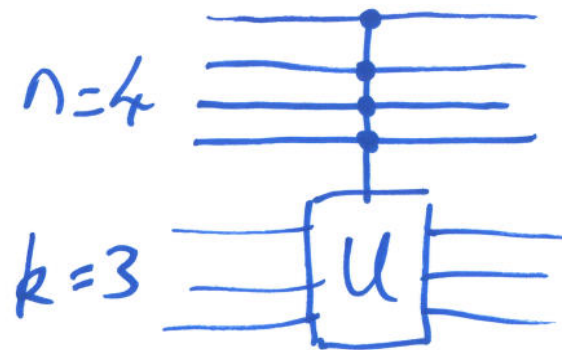


## Multiple Qubit Conditioning

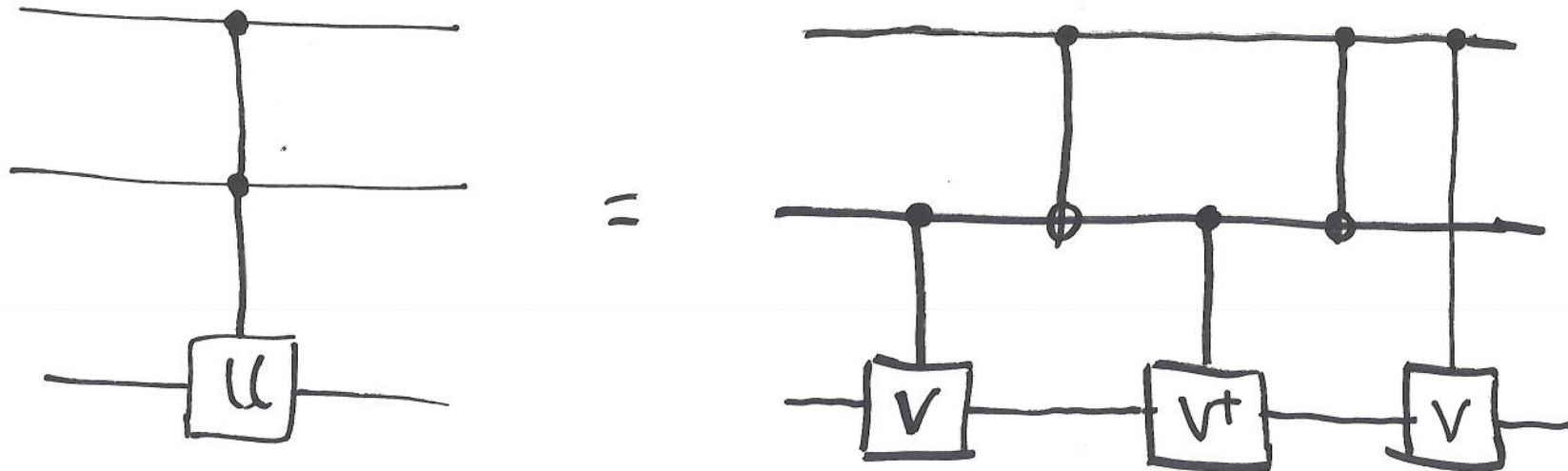
We have  $n+k$  qubits, and  $U$  is a  $k$  qubit unitary.  
Then,

$$\begin{aligned} C^n(U) |x_1 x_2 \dots x_n\rangle |\psi\rangle \\ = |x_1 x_2 \dots x_n\rangle U^{x_1 x_2 \dots x_n} |\psi\rangle \end{aligned}$$

where  $x_1 x_2 \dots x_n$  is the product of bits.



# Multi-Qubit Control via Single-Qubit Control



$C^2(U)$

where  $V^2 = U$ .

When  $V = (1-i)(I+iX)/2$ , then  $C^2(U)$  is the Toffoli gate (i.e.  $V^2 = X$ ).

## Universal Computation

One and two qubit reversible gates are sufficient to implement the Toffoli gate,

... and, more generally,

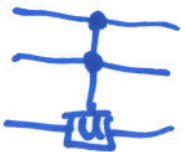
suffice for universal computation.



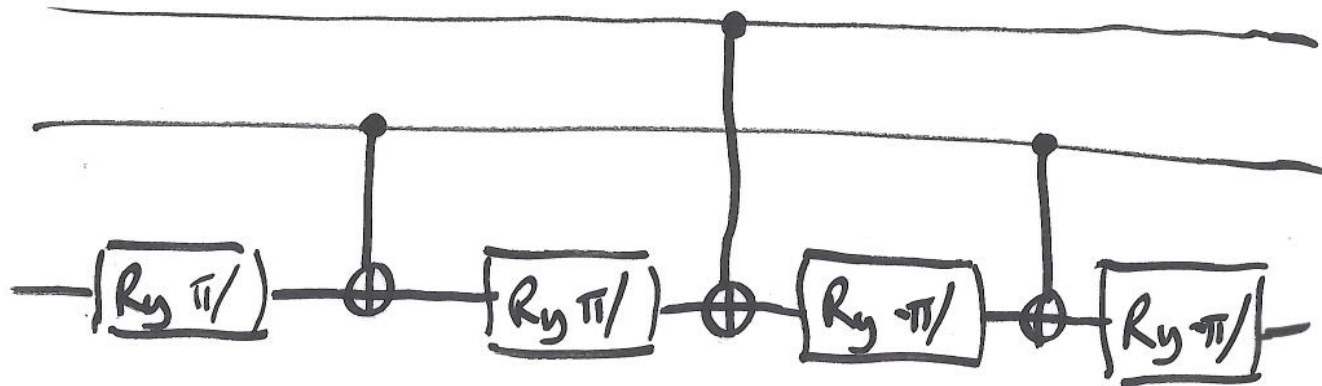


# Fredkin Gate (controlled-swap)

$$\begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix}$$

- 1) Fredkin gate using three Toffoli gates? - (think of swap gate)
- 2) Show that first and last gates can be replaced by CNOT gates.
- 3) Now replace middle Toffoli gate with circuit for  using two qubit gates. (Fig 4.8)
- 4) A simpler construction using only five two-qubit gates?

Observe:



$$|c_1, c_2, t\rangle \rightarrow e^{i\Theta(c_1, c_2, t)} |c_1, c_2, t \oplus c_1 \cdot c_2\rangle$$

$\equiv$  Toffoli up to relative phase

.... Sometimes easier to implement than Toffoli.