Computer Tomography: Computational theory and methods

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Outline

Introduction

- Fields of applications
- Principles of X-ray CT
- Reconstruction techniques in 2D
 - Transform based methods (FBP)
 - Algebraic reconstruction techniques
- Con-beam reconstructions
 - Circular
 - Helical

Conclusions

Definition

Tomography: technique for imaging cross-sections or slices from a set of external measurements of a spatially varying function.

- medical imaging
- industrial nondestructive evaluation
- airport screening
- microtomography

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 - electromagnetic waves

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- CT principles are also applied in:
 - oceanography
 - astronomy
 - geophysics
 - optics

Radon 1917

- reconstruction of a function from its projections

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since then the field is steadily evolving and givis rise to never ceasing flow of new algorithms

Principles of X-ray CT



Principles of X-ray CT

- X-ray traverses an object along stight line
- attenuated signals from various directions are recorded
- reconstruction algorithms are used to convert the projections



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Let

- f(x,y) represents the density of a two-dimensional object
- *I_{in}* applied intensity form the X-ray source
- I_{out} the attenuated intensity of a ray as it propagates throught the object along the line $L = L(\theta, t)$

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 I_{out} the attenuated intensity of a ray as it propagates throught the object along the line $L = L(\theta, t)$

then

$$\int_{L} f(x, y) ds = -\ln \frac{I_{out}}{I_{in}} = P_{\theta}(t)$$

where

- θ projection angle
- *t* detector position

$$P_{\theta}(t) = \int_{L} f(x, y) ds$$



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$$P_{\theta}(t) = \int_{L} f(x, y) ds$$

The line integral can be written as Radon Transform

$$P_{\theta}(t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) \delta(r(x, y, t)) dx dy$$

where

 $\boldsymbol{\delta}$ is the Dirac delta function and

 $r(x, y, t) = \cos \theta + y \sin \theta - t$ (the equation of the X-ray)



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Fourier Slice Theorem

Consider the two-dimensional FT of the object function

$$F(u,v) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y) e^{-j2\pi(ux+vy)} dxdy$$

and the one-dimensional FT of the projection $P_{\theta}(t)$

$$S_{\theta}(w) = \int_{-\infty}^{\infty} P_{\theta}(t) e^{-j2\pi w t} dt$$

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Simplest case, when v = 0 ($\theta = 0$)

$$F(u,0) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x,y)e^{-j2\pi ux} dx dy = \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(x,y) dy \right] e^{-j2\pi ux} dx$$
$$= \int_{-\infty}^{\infty} \left[\int_{L} f(x,y) ds \right] e^{-j2\pi ux} dx = \int_{-\infty}^{\infty} P_{\theta=0}(x)e^{-j2\pi ux} dx = S_{\theta=0}(u)$$

Fourier Slice Theorem

In general for any anglr θ

 $S_{\theta}(w) = F(w, \theta) = F(w \cos \theta, w \sin \theta).$



Fourier reconstruction algorithm

- Take projection of an object function at angles $\theta_1, \theta_2, \ldots, \theta_k$
- Fourier transform each projection to obtain the values of F(u,v)
- Recover the object function f(x, y) by using the inverse FT

$$f(x,y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(u,v) e^{j2\pi(ux+vy)} dudv$$

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Note: infinite number of projections on an interval with lenght π allows exact reconstuction.

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$$f(x,y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(u,v) e^{j2\pi(ux+vy)} dudv$$

Discretising on the square -A/2 < x < A/2; -A/2 < y < A/2

$$f(x,y) \approx \frac{1}{A^2} \sum_{m=-N/2}^{N/2} \sum_{n=-N/2}^{N/2} F\left(\frac{m}{A}, \frac{n}{A}\right) e^{j2\pi((m/A)x + (n/A)y)},$$

where N is even integer.

The values of F(u,v) on a square grid are obtained from the available values along the

radial lines by interpolation.

$$\begin{aligned} f(x,y) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(u,v) e^{j2\pi(ux+vy)} du dv \\ &= \int_{0}^{2\pi} \int_{0}^{\infty} F(w,\theta) e^{j2\pi w} \underbrace{\left(\frac{t}{v \cos \theta} + y \sin \theta \right)}_{w dw d\theta} \end{aligned}$$

$$\begin{split} f(x,y) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(u,v) e^{j2\pi(ux+vy)} du dv \\ &= \int_{0}^{2\pi} \int_{0}^{\infty} F(w,\theta) e^{j2\pi w} \underbrace{\underbrace{}_{v \cos \theta} + y \sin \theta}_{W dw d\theta} \\ &= \int_{0}^{\pi} \underbrace{\int_{-\infty}^{\infty} S_{\theta}(w) |w| e^{j2\pi w t} dw d\theta}_{Q_{\theta}(t)} \end{split}$$

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Step 1: Filtering

$$Q_{\theta}(t) = \int_{-\infty}^{\infty} S_{\theta}(w) |w| e^{j2\pi w t} dw$$

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$$Q_{\theta}(t) = \int_{-\infty}^{\infty} S_{\theta}(w) |w| e^{j2\pi w t} dw$$

Use badlimited FT to compute $S_{\theta}(w)$ i.e.

$$S_{\theta}(w) \approx S_{\theta}\left(m\frac{2W}{N}\right) = \frac{1}{2W} \sum_{k=-N/2}^{N/2-1} P_{\theta}\left(\frac{k}{2W}\right) e^{-j2\pi(mk/N)}, \quad m = -\frac{N}{2}, \dots, \frac{N}{2},$$

where W is a frequency band limit and N is some (large) number corresponding to the number of the sampling points. (FFT)

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where W is a frequency band limit and N is some (large) number corresponding to the number of the sampling points. (FFT) Therefore

$$Q_{\theta}\left(\frac{k}{2W}\right) \approx \left(\frac{2W}{N}\right) \sum_{m=-N/2}^{N/2} S_{\theta}\left(m\frac{2W}{N}\right) \left|m\frac{2W}{N}\right| e^{j2\pi(mk/N)}, \quad k = -\frac{N}{2}, \dots, \frac{N}{2}.$$

Inverse DFT of the product $S_{\theta}(m(2W/N))$ and |m(2W/N)|.

Step 1: Filtering

Superior results are obtained by deemphasizing the high frequencies with the Hamming window function H.

$$Q_{\theta}\left(\frac{k}{2W}\right) \approx \left(\frac{2W}{N}\right) \sum_{m=-N/2}^{N/2} S_{\theta}\left(m\frac{2W}{N}\right) \left|m\frac{2W}{N}\right| H\left(m\frac{2W}{N}\right) e^{j2\pi(mk/N)}$$

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where $\phi(k/"W)$ is the inverse DFT of |m(2W/N)|H(m(2W/N)).

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Step 2: Backprojection Reconstruc f(x, y) by the formula

$$f(x,y) \approx \frac{\pi}{K} \sum_{i=1}^{K} Q_{\theta_i}(x \cos \theta_i + y \sin \theta_i)$$

where K is the number of the projection angles θ_i .

Note The values of Q_{θ} in Step 2 can be approxiamted from the values obtained in Step 1 by interpolation (linear).
The FBP algorithm

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Computational cost: Backprojection summation $\longrightarrow \mathcal{O}(N^3)$

Fast backprojections (Nilsson 96) $\longrightarrow \mathcal{O}(N^2 \log N)$

Equiangular Rays

Let $R_{\beta}(\gamma)$ be a fan projection

- β is the angle between the source and a reference axis.
- γ gives the location of a ray
- D = |SO| source to origin dist.

L - source to point (x, y) dist. In polar coordinates $[(r, \phi)]$

 $L = \sqrt{D + (r \sin \eta)^2 + (r \cos \eta)^2},$ where $\eta = \beta - \phi$.



Equiangular Rays

$$f(x,y) = \int_0^\pi \int_{-\infty}^\infty P_\theta(t)h(x\cos\theta + y\sin\theta - t)dtd\theta$$
$$f(r,\phi) = \frac{1}{2}\int_0^{2\pi} \int_{-\infty}^\infty P_\theta(t)h(r\cos(\theta - \phi) - t)dtd\theta$$

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$$= \dots$$
$$= \int_0^{2\pi} \frac{1}{L^2} Q_{\beta}(\gamma')d\beta,$$

where

 $Q_{\beta}(\gamma) = R'_{\beta}(\gamma) * g(\gamma)$ $R'_{\beta}(\gamma) = R_{\beta}(\gamma) \cdot D \cdot \cos \gamma$ $g(\gamma) = \frac{1}{2} \left(\frac{\gamma}{\sin \gamma}\right)^2 h(\gamma)$

Equiangular Rays

 α is the sampling interval of the projection

 β_i are the angles at which projections are taken

Step 1 For each $R_{\beta_i}(n\alpha), n \in \mathbb{N}$, generate the corresponding modified projection $R'_{\beta_i}(n\alpha) = R_{\beta}(n\alpha) \cdot D \cdot \cos n\alpha$

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Step 2 Convolve $R'_{\beta_i}(n\alpha)$ with $g(n\alpha) = \frac{1}{2} \left(\frac{n\alpha}{\sin n\alpha}\right)^2 h(n\alpha)$ to generate the corresponding filtered projection (FFT)

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 $Q_{\beta_i}(n\alpha) = R'_{\beta_i}(n\alpha) * g(n\alpha) * g(n\alpha),$

where $k(n\alpha)$ is a smoothing filter.

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$$Q_{\beta_i}(n\alpha) = R'_{\beta_i}(n\alpha) * g(n\alpha) * h(n\alpha),$$

where $k(n\alpha)$ is a smoothing filter.

Step 3 Perform a weighted backprojection

$$f(x,y) \approx \Delta\beta \sum_{i=1} K \frac{1}{L^2(x,y,\beta_1)} Q_{\beta_i}(\gamma'),$$

where γ' is the angle of a ray that passes through the point (x, y) and $\Delta\beta = 2\pi/K$.

Equally Spaced Detectors

Let $R_{\beta}(s)$ be a fan projection

- β is the angle between the source and a reference axis.
- *s* is the distance along the detector
- D = |SO| source to origin dist.

U - the dist. source to projection of a pixel on the central ray In polar coordinates $[(r, \phi)]$

$$U = \frac{D + r\sin(\beta - \phi)}{D}$$



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Equally Spaced Detectors

$$f(x,y) = \int_0^{\pi} \int_{-\infty}^{\infty} P_{\theta}(t)h(x\cos\theta + y\sin\theta - t)dtd\theta$$
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$$= \dots$$
$$= \int_0^{2\pi} \frac{1}{U^2} Q_{\beta}(s')d\beta,$$

where

$$Q_{\beta}(s) = R'_{\beta}(s) * g(s)$$
$$R'_{\beta}(s) = R_{\beta}(s) \cdot \frac{D}{\sqrt{D^2 + s^2}}$$
$$g(s) = \frac{1}{2}h(s)$$

Algebraic reconstruction techniques

Advantages

- easy to handle different rays geometry
- more adaptable to missing data better image quality
- metal artifacts are reduced
- less radiation dose

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- easy to handle different rays geometry
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Disadvantage

Higher computational cost!

Image and Projection Representation

The idea: Assume that the cross section consists of array of unknowns

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N - total number of cells $f_j = const$ is the value of f(x, y)in the j^{-th} cell p_i - the i^{-th} line integral *ray-sum*



Image and Projection Representation

The idea: Assume that the cross section consists of array of unknowns

N - total number of cells $f_j = const$ is the value of f(x, y)in the j^{-th} cell p_i - the i^{-th} line integral *ray-sum*

 $\sum_{j=1}^{N} w_{ij} f_j = p_i, \ i = 1, 2, \dots, M$

where M is the total number of rays w_{ij} represents the contribution of the j^{-th} cell to the i^{-th} ray integral



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Therefore we have to solve the following linear system of equations

 $w_{11}f_1 + w_{12}f_2 + \dots + w_{1N}f_N = p_1$ $w_{21}f_1 + w_{22}f_2 + \dots + w_{2N}f_N = p_2$ \vdots $w_{M1}f_1 + w_{M2}f_2 + \dots + w_{MN}f_N = p_M$

for large values of M and N.

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for large values of M and N.

Use iterative methods - projections (Kaczmarz 37, Tanabe 71)

Let
$$f^{(0)} = \left(f_1^{(0)}, f_2^{(0)}, \dots f_N^{(0)}\right)$$
 is the initial guess

$$f^{(i)} = f^{(i-1)} - \frac{\left(f^{(i-1)} \cdot w_i - p_i\right)}{w_i \cdot w_i} w_i \quad i = 1, 2, \dots,$$

where $w_i = (w_{i1}, w_{i,2}, \dots w_{iN})$

Theorem (Tanabe'71)

If there exists a unique solution $f^{(s)}$ to the system of equations, then

 $\lim_{k \to \infty} f^{(kM)} = f^{(s)}$

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The rate of convergence depends of the angles between the hyperplanes

- Full orthogonalization is computationally not feasible
- Pairwise orthogonalization scheme (Ramakrishnani 79)
- Carefully choose the order of the hyperplanes (Hounsield 72)

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The rate of convergence depends of the angles between the hyperplanes

- Full orthogonalization is computationally not feasible
- Pairwise orthogonalization scheme (Ramakrishnan 79)
- Carefully choose the order of the hyperplanes (Hounsfield 72)
- When M > N the process oscillate in the neighborhood of the intersections of the hyperplanes
- When M < N the process converges to a solution $f'^{(s)}$, such that $|f^{(0)} f'^{(s)}|$ is minimized.

Consider again

$$f^{(i)} = f^{(i-1)} - \frac{\left(f^{(i-1)} \cdot w_i - p_i\right)}{w_i \cdot w_i} w_i \quad i = 1, 2, \dots,$$

It can also be written as

$$f_j^{(i)} = f_j^{(i-1)} + \frac{p_i - q_i}{\sum_{k=1}^N w_{ik}^2} w_{ij} \quad j = 1, 2, \dots, N, \quad i = 1, 2, \dots,$$

where

$$q_i = f^{(i-1)} \cdot w_i = \sum_{k=1}^N f_k^{(i-1)} w_{ik}.$$

Therefore

$$f_j^{(i)} = f_j^{(i-1)} + \Delta f_j^{(i)},$$

where

$$\Delta f_j^{(i)} = \frac{p_i - q_i}{\sum_{k=1}^N w_{ik}^2} w_{ij}.$$

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The main computational diffuculty is in the calculation, storage and retrieval of

the weight coefficients w_{ij} .

• Replace w_{ij} by 0 or 1, depending wether the center of the j^{-th} cell is within the i^{-th} ray. Thus we can use

$$\Delta f_j^{(i)} = \frac{p_i - q_i}{N_i}$$

for all cells whos centers are within the $i^{\text{-th}}$ ray. $N_i = \sum_{k=1}^N w_{ik}^2$ is the number of cells whose centers are within the $i^{\text{-th}}$ ray.

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- ART suffer from salt and pepper noise, caused by the approximations used for w_{ij} .
- It is possible to reduce the noce by *relaxation* (i.e. update a pixel by $\lambda_i \Delta f_i^{(i)}$, $\lambda < 1$)

SIRT - change the cell values after going through all of the equations

First reported by Anderson and Kak '84

Reduce the error in the approximation of the ray integrals

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 Reduce the error in the approximation of the ray integrals Insted of solving

$$p_i = \sum_{j=1}^N w_{ij} f_j, \quad i = 1, 2, \dots, M.$$

We use again the Radon Transform

$$p_i = R_i f(x, y) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) \delta(r_i(x, y)) dx dy,$$

where $r_i(x, y) = 0$ is the equation of the *i*^{-th} ray and R_i is the projection operator.

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Let $\{b_j(x, y)\}_{j=1}^N$ be basis functions and let

$$f(x,y) \approx \widetilde{f}(x,y) = \sum_{j=1}^{N} g_j b_j(x,y)$$

then

$$p_i = R_i f(x, y) \approx R_i \widetilde{f}(x, y) = \sum_{j=1}^N g_j a_{ij}$$

wher $a_{ij} = R_i b_j(x, y)$.

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- The two systems are the same for the *pixel basis*

$$b_j(x,y) = \begin{cases} 1 & \text{inside the } j^{-\text{th}} \text{ pixel} \\ 0 & \text{everywhere else} \end{cases}$$

- Superior reconstructions are obtained by using bilinear elements pyramid shaped basis

First reported by Anderson and Kak '84

 Reduce the error in the approximation of the ray integrals bilinear elements insted of the traditional pixel basis

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The overall SART iterative scheme is

$$g_j^{(k+1)} = g_j^{(k)} + \frac{\lambda_k}{\sum_{i=1}^M a_{ij}} \sum_{i=1}^M a_{ij} \frac{p_i - a_i \cdot g^{(k)}}{\sum_{j=1}^N a_{ij}}, \quad j = 1, 2, \dots, N,$$

where a_i is the *i*^{-th} row of A and $\lambda_k > 0$ are relaxation coefficients.

More iterative algorithms

Censor and Elfving'02

The idea: Use generalized projections and diagonal weighting matrices to reflect the sparsity of the A matrix.

Definition: A family $\{G_i\}_{i=1}^M$ of diagonal $N \times N$ matrices with diagonal elements $g_{ij} > 0$ is called *Sparsity Pattern Oriented* (SPO) with respect to $M \times N$ matrix A if:

1) $\sum_{i=1}^{M} \overline{G_i} = I$ 2) $\overline{g_{ij}} = 0$ iff $a_{ij} = 0$ for every i = 1, 2, ..., M.

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Diagonal Weighting (DWE)

$$g_j^{(k+1)} = g_j^{(k)} + \lambda_k \sum_{\substack{i=1\\g_{ij}\neq 0}}^M a_{ij} \frac{p_i - a_i \cdot g^{(k)}}{\sum_{\substack{i=1\\g_{il}\neq 0}}^N \frac{a_{il}}{g_{il}}}, \quad j = 1, 2, \dots, N.$$
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Special case

$$g_{ij} = \begin{cases} \frac{1}{s_j} & \text{if } a_{ij} \neq 0\\ 0 & \text{if } a_{ij} = 0 \end{cases}$$

where s_j is the number of the nonzero elements in column j of A. Component Averaging (CAV)

$$g_j^{(k+1)} = g_j^{(k)} + \lambda_k \sum_{i=1}^M a_{ij} \frac{p_i - a_i \cdot g^{(k)}}{\sum_{l=1}^N s_l a_{il}^2}, \quad j = 1, 2, \dots, N.$$

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• Jiang and Wang 03 ART, SART, CAV, DWE are special cases of the Landweber method $q^{(k+1)} = q^{(k)} + \lambda_k V^{-1} A^* W(p - Aq^{(k)})$

3D Reconstructions

Slice by Slice

Disadvantages

- low speed
- discrete nature
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3D Reconstructions

Slice by Slice

Disadvantages

- low speed
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- Cone-Beam Reconstructions
 - Circular
 - Helical

Circular Cone-Beam Reconstructions



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The Idea: Apply fan-beam FBP to each plane in the cone Feldkemp, Davis, Kress '84 (FDK)



Circular Cone-Beam Reconstructions

The Idea: Apply fan-beam FBP to each plane in the cone Feldkamp, Davis, Kress '84 (FDK)

A ray from the cone is described by

 $t = x \cos \theta + y \sin \theta$ $r = -(-x \sin \theta + y \cos \theta) \sin \gamma$ $+z \cos \gamma$

where

 (t, θ) locate a ray in a fan (r, γ) specify the location of the fan

 θ - projection angle

 γ - fan angle



 $(x, y, z) \longrightarrow (t, s, z)$ (rotation by θ degrees around the *z*-axes)

The FDK Algorithm

Based on the FBP Algorithm for equispatial rays

$$f(t,s,z) = \frac{1}{2} \int_0^{2\pi} \frac{D^2}{(D-s)^2} \int_{-\infty}^{\infty} R_{\beta}(p,\xi) \frac{D}{\sqrt{D^2 + \xi^2 + p^2}} h\left(\frac{Dt}{D-s} - p\right) dp d\beta,$$

where $D = |SO|$ and $\xi = \frac{D}{D-s} z$.

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The Algorithm

Step 1 Compute the modified projection

$$R'_{\beta}(p,\xi) = \frac{D}{\sqrt{D^2 + \xi^2 + p^2}} R_{\beta}(p,\xi)$$

Step 2 Convolve $R'_{eta}(p,\xi)$ with h(p)/2

$$Q_{\beta}(p,\xi) = R'_{\beta}(p,\xi) * \frac{1}{2}h(p)$$

Step 3 Backproject over the 3D reconstruction grid

$$f(t,s,z) = \int_0^{2\pi} \frac{D^2}{(D-s)^2} Q_\beta(\frac{Dt}{D-s}, \frac{Dz}{D-s}) d\beta$$

The Tuy–Smith sufficiency condition

Exact reconstruction is possible if all planes intersecting the object also intersect the source trajectory at least once.

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Exact reconstruction form helical cone-beam data is possible. Tam 95 Kudo, Noo, Defrise 98

Conclusions

The challenges:

- Three dimentional images
- Better quality
- Faster algorithms
- Restricted in time

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